# Analytical Model and Finite-Element Model of an Inductive Displacement Sensor With a Highly Conductive Marker

Dries Doornaert<sup>1</sup>, Christ Glorieux<sup>2</sup>, Herbert De Gersem<sup>1,3</sup>, Robert Puers<sup>4</sup> Werner Spileers<sup>5</sup> and Johan Blanckaert<sup>5,6</sup>

 $1$ KU Leuven - Kulak, Wave Propagation and Signal Processing Research Group, 8500 Kortrijk, Belgium, dries.doornaert@kuleuven-kulak.be

<sup>2</sup>KU Leuven, Laboratory for Acoustics and Thermal Physics, 3001 Heverlee, Belgium

<sup>3</sup>TU Darmstadt, Institut für Theorie Elektromagnetischer Felder (TEMF), 64289 Darmstadt, Germany

<sup>4</sup>KU Leuven, ESAT - MICAS, 3001 Heverlee, Belgium

<sup>5</sup>University Hospitals Leuven, Department of Ophthalmology, 3000 Leuven, Belgium

 ${}^{6}$ Eye & Refractive Center, 8900 Ieper, Belgium

The displacement of a metal sphere causes a change of inductance in a measurement coil. This effect is adequately represented by a finite-element model as well as by a semi-analytical formula, both validated by measurements. The displacement sensor is implanted in an artificial lens for cataract patients.

*Index Terms*—Analytical models, Inductance measurement, Magnetic sensors, Numerical simulation.

# I. INTRODUCTION

THE contraction of the ciliary muscle controls the focus<br>of the eye. An implanted metal spherical marker causes HE contraction of the ciliary muscle controls the focus a change of the inductance of a coil placed in the eye's lens (Fig. [1\)](#page-0-0) [\[1\]](#page-1-0), [\[2\]](#page-1-1), [\[3\]](#page-1-2). The interaction between the highly conductive marker and the coil is highly complicated [\[4\]](#page-1-3), [\[5\]](#page-1-4). This paper develops a cheap model to be implemented in the electronic control scheme.

#### II. APPROXIMATE MODEL

According to the applied frequencies (a few MHz) and the electromagnetic properties of human tissue and metal, wave lengths of several meters are expected in tissue and a skin depth of a few micrometer is expected in the metal sphere. This motivates the use of the magnetoquasistatic approach with the marker modelled as perfectly electric conducting (PEC) boundary [\[6\]](#page-1-5), [\[7\]](#page-1-6). The marker causes a partial shielding of the magnetic field. The impact of the marker is assumed to be proportional to the magnetic energy  $W_d$  expelled by the marker's volume. In an axial configuration (Fig. [2\)](#page-1-7) and using Biot-Savart's law [\[8\]](#page-1-8), the relative change of inductance can be approximated by [\[2\]](#page-1-1)

<span id="page-0-1"></span>
$$
\frac{\Delta L}{L_*} = \frac{\Delta W}{W_*} = \frac{2}{3} \frac{\mathcal{R}^3}{(\mathcal{R}^2 + \delta^2)^3}.
$$
 (1)

Here,  $\mathcal{R} = \frac{R}{r}$  is the coil radius and  $\delta = \frac{d}{r}$  is the displacement, both relative to the marker radius,  $L_*$  and  $W_*$  are the inductance and the magnetic energy in the situation without marker and  $\Delta L$  and  $\Delta W = W_* - W_d$  are the changes thereof due to the marker's presence. Model [1](#page-0-1) is sufficiently cheap and allows inclusion in a simple electronic control scheme, but may lack accuracy due to the severe assumptions involved.



<span id="page-0-0"></span>Fig. 1. Marker implanted in the ciliary muscle at a distance between 2 and 4 mm. When the eye accommodates the ciliary muscle moves the marker closer or further away from the sensor (maximum range of 1 mm), causing a change of inductance of the coil [\[2\]](#page-1-1).

### III. FINITE-ELEMENT MODEL

The approximate model is validated by a set of 2D axisymmetric finite-element (FE) models (Fig. [2\)](#page-1-7). The marker is modelled as a PEC boundary. For the considered frequencies, the relative permittivity and permeability of human tissue can be considered as 1 [\[9\]](#page-1-9). The coil and sphere are embedded in a sufficiently large computational domain to account for stray fluxes [\[10\]](#page-1-10). The FE results are in good agreement with the results of the approximate model (Fig. [3\)](#page-1-11). Remarkable is the fact that also at larger distances  $(d > 3$  mm), both approximate model and the FE model feature a sufficient resolution.

The explicit change of the marker's position necessitate remeshing between two successive FE runs. An alternative approach involves a conformal mapping of a cylindrical region between coil and marker with varying height  $H(d)$  to a reference cylinder with fixed height  $H_{\text{ref}} = H(d_{\text{ref}})$ . Then, instead of changing the model geometry and remeshing, the displacement is modelled by a varying, fictitious, anisotropic permeability  $(\tilde{\mu}_r, \tilde{\mu}_z) = (\xi \mu_0, \mu_0/\xi)$  with  $\xi = d/d_{\text{ref}}$ , which is assigned to the reference cylinder region.



<span id="page-1-7"></span>Fig. 2. Axially symmetric 2D model of a coil with  $h = 0.42$  mm, thickness  $t = 0.21$  mm, radius  $R = 1$  mm and number of turns  $N = 1$  enclosed in a fixed modelling sphere with radius  $k = 20$  mm together with a spherical marker with radius  $r$  placed at a distance  $d$  from the face of the coil. Enlarging the coil radius by a factor also enlarges  $t,h$  and  $k$  by the same factor.



<span id="page-1-11"></span>Fig. 3. Simulated data with  $\mathcal{R} = 3$  as function of marker distance d plotted together with Eq. [1](#page-0-1) for the corresponding coil radius  $R$  and marker radius  $r$ .

# IV. VALIDATION BY MEASUREMENT

An experiment was set up with aluminium spheres of different diameter (Fig. [4\)](#page-1-12). The inductance change is measured as a shift of the resonance frequency of a Colpitts oscillator, which is observed as a change of DC voltage in a phase lock loop [\[11\]](#page-1-13). The expected frequency shift is

<span id="page-1-14"></span>
$$
\Delta f = f_* - f_d \simeq -\frac{1}{3} \frac{\mathcal{R}^3}{(\mathcal{R}^2 + \delta^2)^3} f_* \tag{2}
$$

with  $f_*$  the frequency without sphere and  $f_d$  the frequency with a sphere at distance  $d$  [\[2\]](#page-1-1). The measured frequency shift corresponds very well to the frequency shift predicted by [\(2\)](#page-1-14) (Fig. [5\)](#page-1-15).

### V. CONCLUSIONS

An inductive distance measurement system is equipped with a model that relates a change in resonance frequency to the distance between a metal sphere and a measurement coil. The validation by a FE model and verification by measurements show a very good agreement. The technique is mature for inclusion in a artificial lens with accommodation.

# VI. ACKNOWLEDGEMENT

The work of Dries Doornaert is funded by the Research



<span id="page-1-12"></span>Fig. 4. Experimental setup consisting of a spherical marker placed on a wooden stick, a stepper motor, Colpitts oscillator with detection coil, phase locked loop, oscilloscope and a pc (oscilloscope and pc not depicted here). The stepper motor moves the spherical marker along the central axis of the coil thereby changing the inductance  $L$  of the coil.



<span id="page-1-15"></span>Fig. 5. Experimentally measured frequency shift for  $\mathcal{R} = 2$  as function of scaled distance  $\delta$  plotted together with Eq. [2](#page-1-14) (full line).

#### **REFERENCES**

- <span id="page-1-0"></span>[1] D. Doornaert, C. Glorieux, H. De Gersem, R. Puers, W. Spileers, and J. Blanckaert, "Intraocular electro-optic lens with ciliary muscle controlled accommodation," *Annual International Conference of the IEEE Engineering in Medicine and Biology Society*, pp. 3190–3, July 2013.
- <span id="page-1-1"></span>[2] D. Doornaert, C. Glorieux, R. Puers, H. De Gersem, W. Spileers, and J. Blanckaert, "Physiological constraints for an intraocular inductive distance sensor," in *36th Annual International Conference of the IEEE Engineering in Medicine and Biology Society (EMBC)*, Chicago, USA, Aug. 2014, pp. 646–649.
- <span id="page-1-2"></span>[3] C. Glorieux, J. Blanckaert, and R. Puers, "Bionic eye lens PCT/BE2011/000045," 2006.
- <span id="page-1-3"></span>[4] A.A. Kolyshkin and R. Vaillancourt, "Impedance of a single-turn coil due to a double-layered sphere with varying properties," *IEEE Transactions on Magnetics*, vol. 31, no. 3, pp. 2274–2279, May 1995.
- <span id="page-1-4"></span>[5] S. Babic, S. Salon, and C. Akyel, "The mutual inductance of two thin coaxial disk coils in air," *IEEE Transactions on Magnetics*, vol. 40, no. 2, pp. 822–825, Mar. 2004.
- <span id="page-1-5"></span>[6] H.K. Dirks, "Quasi-stationary fields for microelectronic applications," *Electrical Engineering*, vol. 79, pp. 145–155, 1996.
- <span id="page-1-6"></span>[7] R.L. Stoll, *The Analysis of Eddy Currents*, Clarendon Press, 1974.
- <span id="page-1-8"></span>[8] F.W. Grover, *Inductance Calculations*, Dover, 1964.
- <span id="page-1-9"></span>[9] S. Gabriel, R.W. Lau, and C. Gabriel, "The dielectric properties of biological tissues: II. Measurements in the frequency range 10 Hz to 20 GHz," *Physics in Medicine and Biology*, vol. 41, pp. 2251–2269, 1996.
- <span id="page-1-10"></span>[10] Q. Chen and A. Konrad, "A review of finite element open boundary techniques for static and quasistatic electromagnetic field problems, *IEEE Transactions on Magnetics*, vol. 33, no. 1, pp. 663–676, Jan. 1997.
- <span id="page-1-13"></span>[11] Emad Hegazi, Jacob Rael, and Asad Abidi, *The Designer's Guide to High-Purity Oscillators*, Kluwer Academic Publishers, 2004.